

Chiral Symmetry and Electromagnetic Probes

Hendrik van Hees

Texas A&M University

April 29, 2005



Outline

QCD and Chiral Symmetry

Vector Mesons and electromagnetic Probes

Challenges for experiment (and theory)

QCD and (“accidental”) Symmetries

- ▶ Theory for strong interactions: QCD

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4} F_a^{\mu\nu} F_{\mu\nu}^a + \bar{\psi} (\mathrm{i} \not{D} - \hat{M}) \psi$$

- ▶ Particle content:

- ▶ ψ : Quarks, including flavor- and color degrees of freedom,
 $\hat{M} = \text{diag}(m_u, m_d, m_s, \dots)$ = current quark masses
- ▶ A_μ^a : gluons, gauge bosons of $SU(3)_{\text{color}}$

QCD and (“accidental”) Symmetries

- ▶ Theory for strong interactions: QCD

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4} F_a^{\mu\nu} F_{\mu\nu}^a + \bar{\psi} (\mathrm{i} \not{D} - \hat{M}) \psi$$

- ▶ Particle content:
 - ▶ ψ : Quarks, including flavor- and color degrees of freedom,
 $\hat{M} = \text{diag}(m_u, m_d, m_s, \dots)$ = current quark masses
 - ▶ A_μ^a : gluons, gauge bosons of $SU(3)_{\text{color}}$
- ▶ Symmetries
 - ▶ fundamental building block: local $SU(3)_{\text{color}}$ symmetry
 - ▶ in light-quark sector: approximate chiral symmetry
 - ▶ dilation symmetry (scale invariance)

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)
- ▶ chiral limit: $\hat{M} \rightarrow 0 \Rightarrow$, vector-axial-vector symmetries
 - ▶ $\psi \rightarrow \exp[-i(\vec{\alpha}_V + \gamma_5 \vec{\alpha}_A) \vec{T}] \psi$
 \vec{T} : generators of $SU(2)_{\text{flavor}}$ (or $SU(3)_{\text{flavor}}$)

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)
- ▶ chiral limit: $\hat{M} \rightarrow 0 \Rightarrow$, vector-axial-vector symmetries
 - ▶ $\psi \rightarrow \exp[-i(\vec{\alpha}_V + \gamma_5 \vec{\alpha}_A) \vec{T}] \psi$
 \vec{T} : generators of $SU(2)_{\text{flavor}}$ (or $SU(3)_{\text{flavor}}$)
 - ▶ **conserved vector and axial-vector currents:**

$$\vec{j}_V^\mu = \bar{\psi} \vec{T} \gamma^\mu \psi, \quad \vec{j}_A^\mu = \bar{\psi} \vec{T} \gamma_5 \gamma^\mu \psi$$

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)
- ▶ chiral limit: $\hat{M} \rightarrow 0 \Rightarrow$, vector-axial-vector symmetries
 - ▶ $\psi \rightarrow \exp[-i(\vec{\alpha}_V + \gamma_5 \vec{\alpha}_A) \vec{T}] \psi$
 \vec{T} : generators of $SU(2)_{\text{flavor}}$ (or $SU(3)_{\text{flavor}}$)
 - ▶ **conserved vector and axial-vector currents:**

$$\vec{j}_V^\mu = \bar{\psi} \vec{T} \gamma^\mu \psi, \quad \vec{j}_A^\mu = \bar{\psi} \vec{T} \gamma_5 \gamma^\mu \psi$$

- ▶ in vacuum, chiral symmetry **spontaneously broken** by quark condensate: $\langle 0 | \bar{\psi} \psi | 0 \rangle \neq 0$

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)

- ▶ chiral limit: $\hat{M} \rightarrow 0 \Rightarrow$, vector-axial-vector symmetries

- ▶ $\psi \rightarrow \exp[-i(\vec{\alpha}_V + \gamma_5 \vec{\alpha}_A) \vec{T}] \psi$

- ▶ \vec{T} : generators of $SU(2)_{\text{flavor}}$ (or $SU(3)_{\text{flavor}}$)

- ▶ **conserved vector and axial-vector currents:**

$$\vec{j}_V^\mu = \bar{\psi} \vec{T} \gamma^\mu \psi, \quad \vec{j}_A^\mu = \bar{\psi} \vec{T} \gamma_5 \gamma^\mu \psi$$

- ▶ in vacuum, chiral symmetry **spontaneously broken** by quark condensate: $\langle 0 | \bar{\psi} \psi | 0 \rangle \neq 0$

- ▶ (approximate) Goldstone bosons: π, K, η (pseudoscalar octet)

"Fate" of chiral symmetry

- ▶ in classical field theory: each continuous symmetry defines **conserved current** (**Noether's theorem**)

- ▶ chiral limit: $\hat{M} \rightarrow 0 \Rightarrow$, vector-axial-vector symmetries

- ▶ $\psi \rightarrow \exp[-i(\vec{\alpha}_V + \gamma_5 \vec{\alpha}_A) \vec{T}] \psi$

\vec{T} : generators of $SU(2)_{\text{flavor}}$ (or $SU(3)_{\text{flavor}}$)

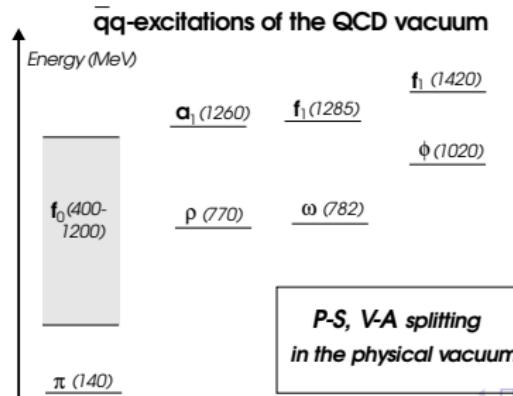
- ▶ **conserved vector and axial-vector currents:**

$$\vec{j}_V^\mu = \bar{\psi} \vec{T} \gamma^\mu \psi, \quad \vec{j}_A^\mu = \bar{\psi} \vec{T} \gamma_5 \gamma^\mu \psi$$

- ▶ in vacuum, chiral symmetry **spontaneously broken** by quark condensate: $\langle 0 | \bar{\psi} \psi | 0 \rangle \neq 0$
- ▶ (approximate) Goldstone bosons: π, K, η (pseudoscalar octet)
- ▶ "real world": chiral symmetry slightly **explicitly broken** by quark masses $\hat{M} \neq 0$

Phenomenology from Chiral Symmetry

- ▶ Use (approximate) **chiral symmetry** to build effective models
- ▶ Ward identities
 - ▶ PCAC: $\langle 0 | \partial^\mu j_{A\mu}^k | \pi^j(\vec{k}) \rangle = iF_\pi m_\pi^2 \delta^{kj}$
 - ▶ $m_\pi^2 F_\pi^2 = -(m_u + m_d) \langle 0 | \bar{u}u | 0 \rangle$ (GOR relation)
- ▶ Spontaneous breaking causes splitting of chiral partners:



Finite Temperature

- ▶ Asymptotic freedom \Rightarrow quark condensate melts at high enough temperatures
- ▶ all bulk properties from partition sum:

$$\mathcal{Z}(V, T, \mu_q) = \text{Tr}\{\exp[-(\mathbf{H} - \mu_q \mathbf{N})/T]\}$$

Finite Temperature

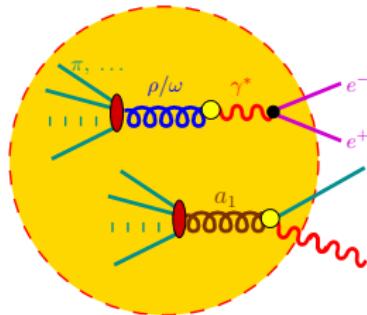
- ▶ Asymptotic freedom \Rightarrow quark condensate melts at high enough temperatures
- ▶ all bulk properties from partition sum:

$$Z(V, T, \mu_q) = \text{Tr}\{\exp[-(\mathbf{H} - \mu_q \mathbf{N})/T]\}$$

- ▶ Free energy: $\Omega = -\frac{T}{V} \ln Z = -P$
- ▶ Quark condensate: $\langle \bar{\psi}_q \psi_q \rangle_{T, \mu_q} = \frac{V}{T} \frac{\partial P}{\partial m_q}$
- ▶ Lattice QCD indicates: Chiral symmetry restoration \leftrightarrow deconfinement phase transition (same T_c)

Why Electromagnetic Probes?

- ▶ γ, ℓ^\pm : only e. m. interactions
- ▶ reflect whole “history” of collision
- ▶ chance to see chiral symm. rest. directly?



Why Electromagnetic Probes?

- ▶ γ, ℓ^\pm : only e. m. interactions
- ▶ reflect whole “history” of collision
- ▶ chance to see chiral symm. rest. directly?

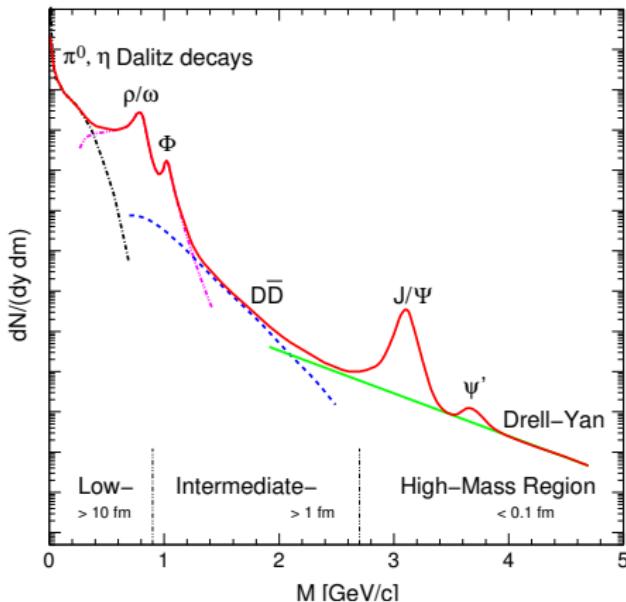
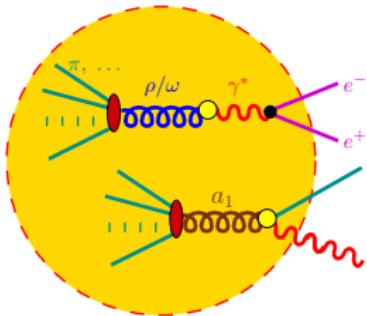


Fig. by A. Drees (from [RW00])

Vector Mesons and electromagnetic Probes

- ▶ photon and dilepton thermal emission rates given by same electromagnetic-current-correlation function

$$(J_\mu = \sum_f Q_f \bar{\psi}_f \gamma_\mu \psi_f)$$

$$\Pi_{\mu\nu}^{<}(q) = \int d^4x \exp(iq \cdot x) \langle J_\mu(0) J_\nu(x) \rangle_T = -2n_B(q_0) \operatorname{Im} \Pi_{\mu\nu}^{(\text{ret})}(q)$$

$$q_0 \frac{dN_\gamma}{d^4x d^3\vec{q}} = \frac{\alpha_{\text{em}}}{2\pi^2} g^{\mu\nu} \operatorname{Im} \Pi_{\mu\nu}^{(\text{ret})}(q) \Big|_{q_0=|\vec{q}|}$$

$$\frac{dN_{e^+e^-}}{d^4x d^4k} = -g_{\mu\nu} \frac{\alpha^2}{3q^2\pi^3} \operatorname{Im} \Pi_{\mu\nu}^{(\text{ret})}(q) \Big|_{q^2=M_{e^+e^-}^2}$$

- ▶ to lowest order in α : $e^2 \Pi_{\mu\nu} \simeq \Sigma_{\mu\nu}^{(\gamma)}$
- ▶ derivable from underlying thermodynamic potential Ω !

Vector Mesons and chiral symmetry

- ▶ **vector** and **axial-vector** mesons \leftrightarrow correlators of the respective currents

$$\Pi_{V/A}^{\mu\nu}(p) := \int d^4x \exp(ipx) \left\langle J_{V/A}^\nu(0) J_{V/A}^\mu(x) \right\rangle_{\text{ret}}$$

Vector Mesons and chiral symmetry

- ▶ vector and axial-vector mesons \leftrightarrow correlators of the respective currents

$$\Pi_{V/A}^{\mu\nu}(p) := \int d^4x \exp(ipx) \left\langle J_{V/A}^\nu(0) J_{V/A}^\mu(x) \right\rangle_{\text{ret}}$$

- ▶ Ward-Takahashi Identities from chiral symmetry \Rightarrow Weinberg-sum rules

$$f_\pi^2 = - \int_0^\infty \frac{dp_0^2}{\pi p_0^2} [\text{Im } \Pi_V(p_0, 0) - \text{Im } \Pi_A(p_0, 0)]$$

Vector Mesons and chiral symmetry

- ▶ vector and axial-vector mesons \leftrightarrow correlators of the respective currents

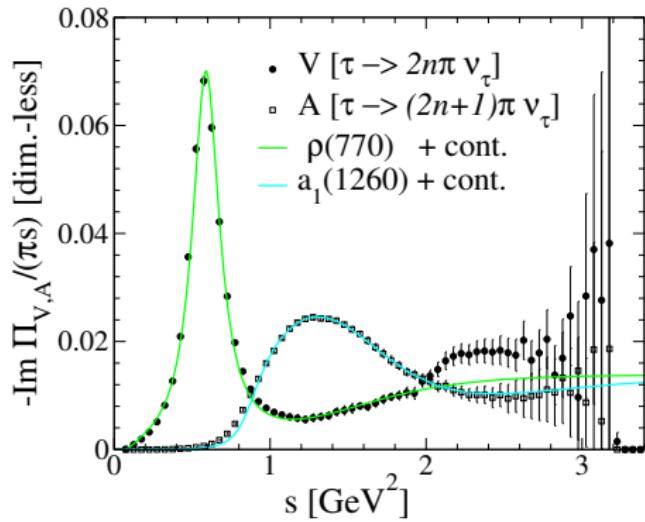
$$\Pi_{V/A}^{\mu\nu}(p) := \int d^4x \exp(ipx) \left\langle J_{V/A}^\nu(0) J_{V/A}^\mu(x) \right\rangle_{\text{ret}}$$

- ▶ Ward-Takahashi Identities from chiral symmetry \Rightarrow
Weinberg-sum rules

$$f_\pi^2 = - \int_0^\infty \frac{dp_0^2}{\pi p_0^2} [\text{Im } \Pi_V(p_0, 0) - \text{Im } \Pi_A(p_0, 0)]$$

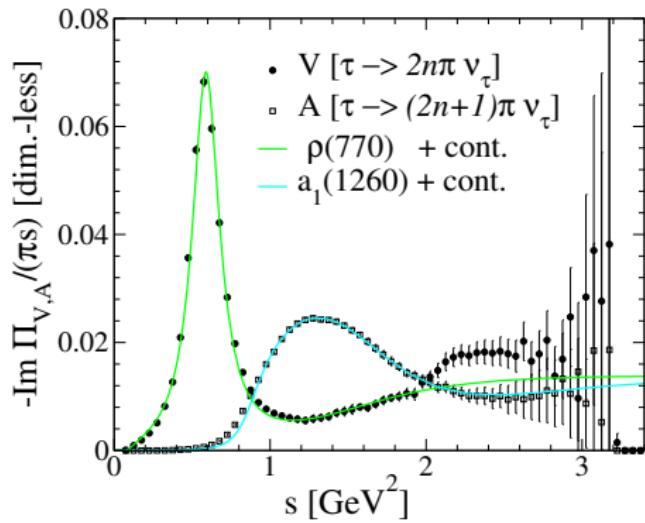
- ▶ spectral functions of vector (e.g. ρ) and axial vector (e.g. a_1) directly related to order parameter of chiral symmetry!

Vector Mesons and chiral symmetry

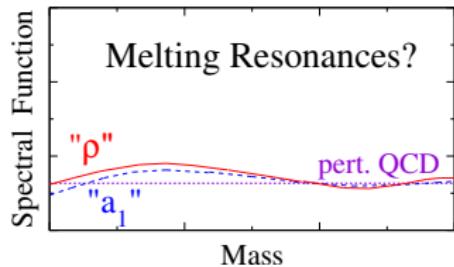
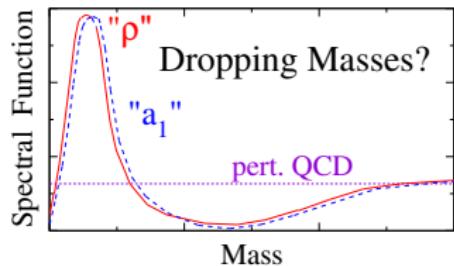


from [Rap03]

Vector Mesons and chiral symmetry



from [Rap03]



from [Rap05]

Models

- ▶ different models with chiral symmetry: equivalent only on shell (**“low-energy theorems”**)

Models

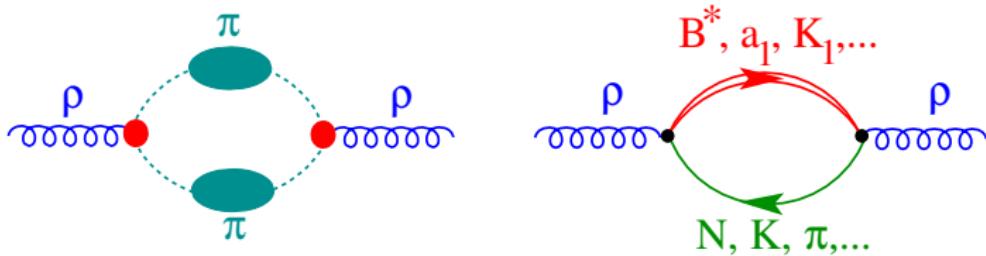
- ▶ different models with chiral symmetry: equivalent only on shell (**“low-energy theorems”**)
- ▶ model-independent conclusions only in **low-temperature/density limit** (chiral perturbation theory) or from **lattice-QCD calculations**

Models

- ▶ different models with chiral symmetry: equivalent only on shell (**“low-energy theorems”**)
- ▶ model-independent conclusions only in **low-temperature/density limit** (chiral perturbation theory) or from **lattice-QCD calculations**
- ▶ use **phenomenological hadronic models** + many-body techniques to assess medium modifications of vector mesons

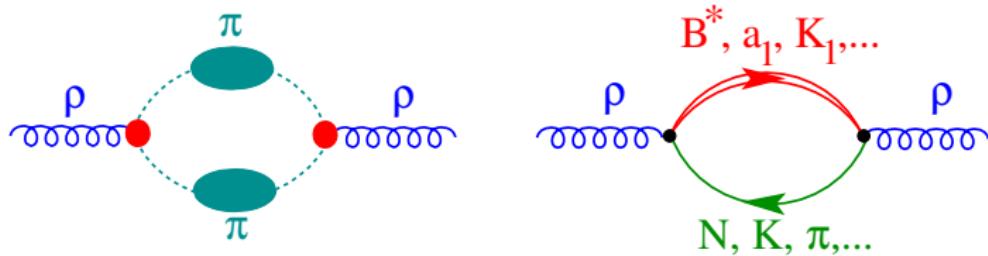
Models

- ▶ Phenomenological hadronic models [Chanfray et al, Herrmann et al, Rapp et al, ...] for vector mesons
- ▶ important ingredients: $\pi\pi$ interactions
baryonic excitations



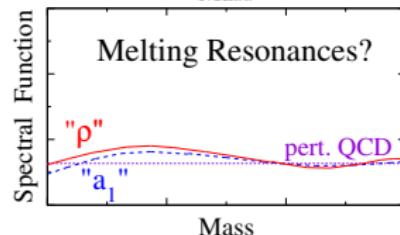
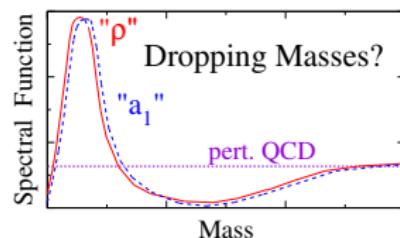
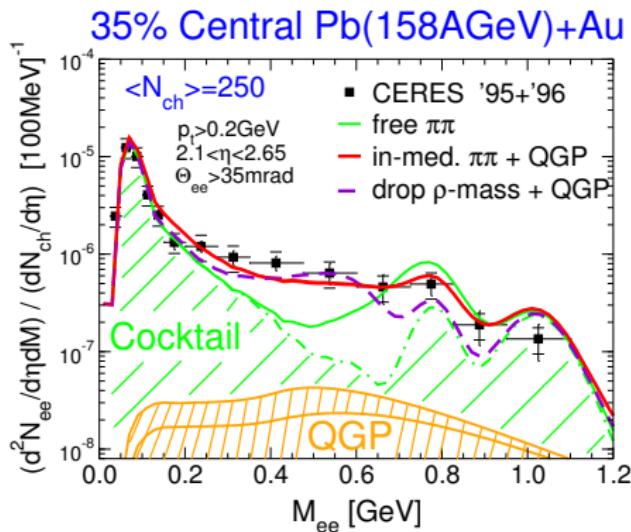
Models

- ▶ Phenomenological hadronic models [Chanfray et al, Herrmann et al, Rapp et al, ...] for vector mesons
- ▶ important ingredients: $\pi\pi$ interactions
baryonic excitations



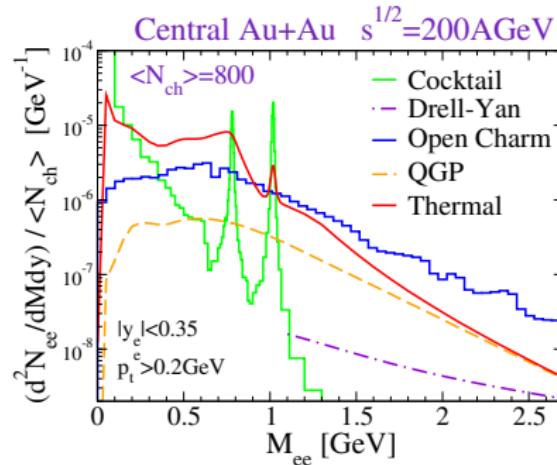
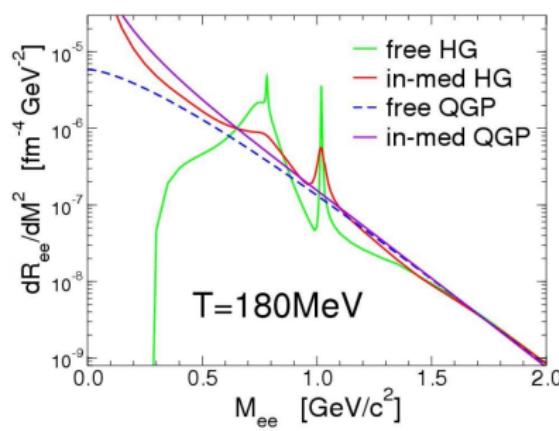
- ▶ **Baryon (resonances)** important, even at RHIC with low **net** baryon density $n_B - n_{\bar{B}}$
- ▶ reason: $n_B + n_{\bar{B}}$ relevant (CP invariance of strong interactions)

Dilepton rates at SpS



- ▶ how to decide about scenario **experimentally?**

Dilepton rates/spectra at RHIC



- ▶ in-medium hadron gas **matches with in-medium pQCD**
- ▶ (similar results also for γ rates)
- ▶ “quark-hadron duality”?

Challenges for Experiment

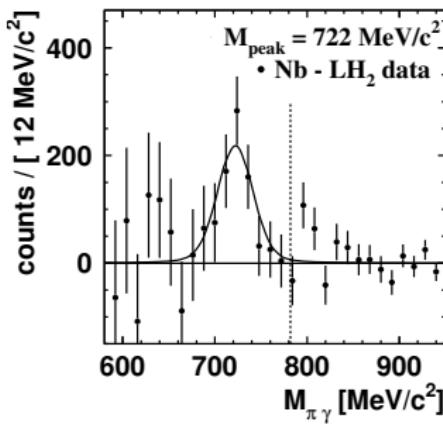
- ▶ Direct signature for chiral restoration:
spectra for ρ and a_1 mesons degenerate
- ▶ $\pi^\pm\gamma$ invariant mass spectrum $\leftrightarrow a_1$ spectral function

X	$\Gamma_{X \rightarrow \pi\gamma} [\text{MeV}]$
a_1	0.64
ρ	0.07
ω	only $\pi^0\gamma$!
a_2	0.3
$\pi(1300)$???

Challenges for Experiment

- ▶ Direct signature for chiral restoration:
spectra for ρ and a_1 mesons degenerate
- ▶ $\pi^\pm\gamma$ invariant mass spectrum $\leftrightarrow a_1$ spectral function

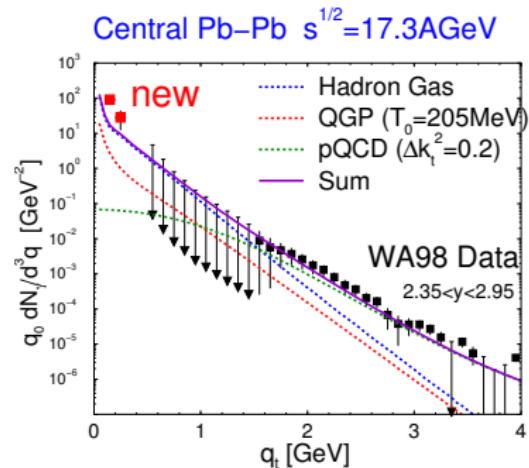
X	$\Gamma_{X \rightarrow \pi\gamma} [\text{MeV}]$
a_1	0.64
ρ	0.07
ω	only $\pi^0\gamma!$
a_2	0.3
$\pi(1300)$???



ω -spectral function from [Trn05] (CBELSA/TAPS)

Challenges for Experiment

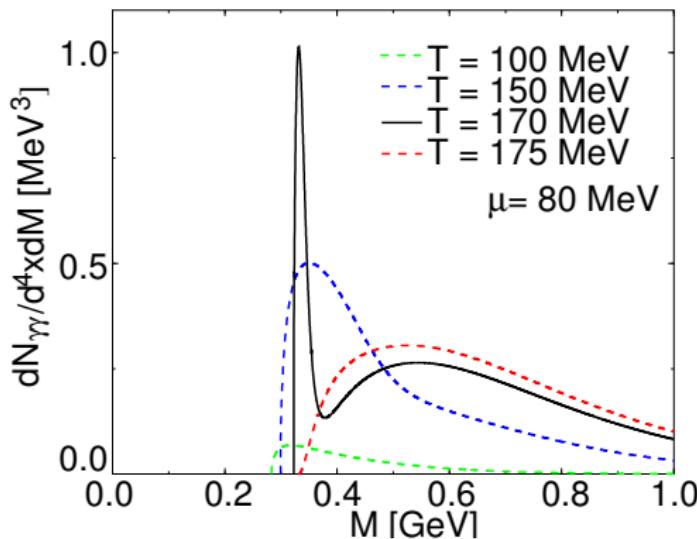
- ▶ Photon rate



- ▶ $\pi\pi \rightarrow \rho \rightarrow \pi\pi\gamma$ not enough to explain enhancement
- ▶ New development (Liu/Rapp work in progress):
 $\pi K \rightarrow K^* \rightarrow \pi K\gamma$

Challenges for Experiment

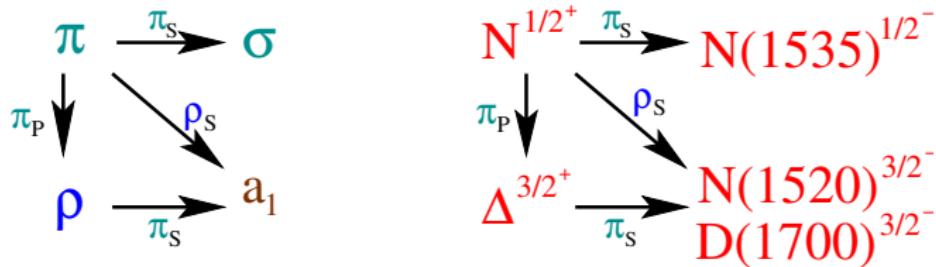
- ▶ low-mass $\gamma\gamma$ -invariant-mass spectrum
- ▶ dropping and narrowing of $\sigma(600)$ for temperatures around T_c



$\pi^+ + \pi^- \rightarrow \gamma + \gamma$ from [VKB⁺98]

Challenges for Theory

- ▶ Need a fully **chiral** model



- ▶ How to treat **(axial-) vector mesons** (gauge model?)
- ▶ Approximation scheme for both **dynamical properties** (spectral functions) and **thermodynamic bulk properties** (phase diagram)?

Summary

- ▶ chiral symmetry: important feature of low-energy sector of QCD
- ▶ one aspect of (s)QGP: how is chiral symmetry restored?
- ▶ electromagnetic probes may provide most direct insight
 - ▶ invariant-mass spectra for chiral partners: here ρ vs. a_1 (“dropping mass” vs. “in-medium broadening”?)
 - ▶ low-energy photons \leftrightarrow dileptons (puzzle?)
 - ▶ $\gamma\gamma$ spectra: σ “softening”?
- ▶ a lot to do also for theory: consistent chiral scheme for hadrons wanted!

Bibliography I

- [Rap03] R. Rapp, Dileptons in high-energy heavy-ion collisions, *Pramana* **60** (2003) 675
- [Rap05] R. Rapp, The vector probe in heavy-ion reactions, *J. Phys.* **G31** (2005) S217, URL
<http://arxiv.org/abs/nucl-th/0409054>
- [RW00] R. Rapp, J. Wambach, Chiral symmetry restoration and dileptons in relativistic heavy-ion collisions, *Adv. Nucl. Phys.* **25** (2000) 1, URL
<http://arXiv.org/abs/hep-ph/9909229>

Bibliography II

- [Trn05] D. Trnka (CBELSA/TAPS), First observation of in-medium modifications of the omega meson (2005), URL <http://arxiv.org/abs/nucl-ex/0504010>
- [VKB⁺⁹⁸] M. K. Volkov, E. A. Kuraev, D. Blaschke, G. Ropke, S. M. Schmidt, Excess low energy photon pairs from pion annihilation at the chiral phase transition (1998), URL <http://arxiv.org/abs/hep-ph/9706350>